### **Continuum Mechanics and the Finite Element Method**



#### **Assignment 2**

#### ◆ Due on March 2<sup>nd</sup> @ midnight

## Suppose you want to simulate this...



## The familiar mass-spring system



Spring length before/after  $\mathbf{l} = |\mathbf{x} - \mathbf{y}_i|$  $\mathbf{l}_0 = |\mathbf{X} - \mathbf{y}_i|$ 

Deformation Measure Elastic Energy Forces  $e = \left(\frac{l}{l_0} - 1\right)$   $W = \frac{1}{2}ke^2$   $f_{int} = -\frac{\partial W}{\partial x}$ 

$$\mathbf{f}_{\text{int}}(\mathbf{x}) = -k \left( \left( \frac{l}{l_0} - 1 \right) \frac{\mathbf{x} - \mathbf{y}_i}{|\mathbf{x} - \mathbf{y}_i|} \right)$$

### Mass Spring Systems

- Can be used to model arbitrary elastic/plastic objects, but...
  - Behavior depends on tessellation
    - Find good spring layout
    - Find good spring constants
    - Different types of springs interfere
    - No direct map to measurable material properties



#### Alternative...

- Start from continuum mechanics
- Discretize with Finite Elements
  - Decompose model into simple elements
  - Setup & solve system of algebraic equations
- Advantages
  - Accurate and controllable material behavior
  - Largely independent of mesh structure





#### Mass Spring vs Continuum Mechanics

Mass spring systems require:

- **1.** Measure of Deformation  $\left(\frac{l}{l_0}-1\right)$
- **2.** Material Model k
- 3. Deformation Energy W
- 4. Internal Forces  $f_{int} = -\frac{\partial W}{\partial x}$

$$\frac{W}{\partial W} = \frac{1}{2}ke^2$$

 We need to derive the same types of concepts using continuum mechanics principles

#### Continuum Mechanics: 3D Deformations

- For a deformable body, identify:
  - undeformed state  $\Omega \subset \mathbf{R}^3$  described by positions X - deformed state  $\Omega' \subset \mathbf{R}^3$  described by positions **x**
- Displacement field  ${\bf u}$  describes  $\Omega'$  in terms of  $\Omega$

$$\mathbf{u}: \Omega \to \Omega' \qquad \mathbf{x} = X + \mathbf{u}(X)$$



#### Continuum Mechanics: 3D Deformations

- Consider material points  $X_1$  and  $X_2$  such that  $|\mathbf{d}|$  is infinitesimal, where  $\mathbf{d} = X_2 X_1$
- Now consider deformed vector  $\mathbf{d}'$

d

Deformation gradient 
$$\mathbf{F} = \frac{\partial f}{\partial x}$$
  
 $\mathbf{F} = \mathbf{x}_2 - \mathbf{x}_1 \approx (\mathbf{I} + \nabla \mathbf{u}) d$ 



#### So...

Displacement field transforms points

 $\mathbf{x} = X + \mathbf{u}(X)$ 

 Jacobian of displacement field (deformation gradient) transforms differentials (infinitesimal vectors) from undeformed to deformed

$$\mathbf{d'} = (\mathbf{I} + \nabla \mathbf{u}) \mathbf{d} = \mathbf{F} \mathbf{d} \qquad \mathbf{F} = \frac{\partial \mathbf{X}}{\partial X}$$





$$egin{array}{ll} x &= X\cos heta &- Y\sin heta \ y &= X\sin heta &+ Y\cos heta \ \end{array}$$

$$\mathbf{F} = egin{bmatrix} \cos heta & -\sin heta \ \sin heta & \cos heta \end{bmatrix}$$







#### **Measure of deformations**

• Displacement field transforms points

$$\mathbf{x} = X + \mathbf{u}(X)$$

Jacobian of displacement field (deformation gradient) transforms vectors

$$\mathbf{d'} = (\mathbf{I} + \nabla \mathbf{u}) \mathbf{d} = \mathbf{F} \mathbf{d} \qquad \mathbf{F} = \frac{\partial \mathbf{x}}{\partial X}$$

How can we describe deformations?

### **Back to spring deformation**

◆ Deformation measure (strain):  $\left(\frac{l}{l_0} - 1\right)$ ◆ Undeformed spring:  $\frac{l}{l_0} = 1$ 

Undeformed (infinitesimal) continuum volume:

$$\mathbf{F} = \mathbf{I}$$
?

**Strain (**description of deformation in terms of *relative* displacement**)** 

• Deformation measure (strain):  $\left(\frac{l}{l_0}-1\right)$ 

- Desirable property: if spring is undeformed, strain is 0 (no change in shape)
- Can we find a similar measure that would work for infinitesimal volumes?

#### **3D Nonlinear Strain**

Idea: to quantify change in shape, measure change in squared length for any arbitrary vector

$$|\mathbf{d}'|^2 - |\mathbf{d}|^2 = \mathbf{d}'^T \mathbf{d}' - \mathbf{d}^T \mathbf{d} = \mathbf{d}^T (\mathbf{F}^T \mathbf{F} - \mathbf{I}) \mathbf{d}$$



**Strain (**description of deformation in terms of *relative* displacement**)** 

• Deformation measure (strain):  $\left(\frac{l}{l_0}-1\right)$ 

 Desirable property: if spring is undeformed, strain is 0 (no change in shape)

Can we find a similar measure that would work for infinitesimal volumes?

**Green strain** 
$$\mathbf{E} = \frac{1}{2} (\mathbf{F}^T \mathbf{F} - \mathbf{I})$$

#### **3D Linear Strain**

- Green strain is quadratic in displacements  $\mathbf{E} = \frac{1}{2} (\mathbf{F}^T \mathbf{F} - \mathbf{I}) = \frac{1}{2} (\nabla \mathbf{u} + \nabla \mathbf{u}^T + \nabla \mathbf{u}^T \nabla \mathbf{u})$
- Neglecting quadratic term (small deformation assumption) leads to the linear

**Cauchy strain** (small strain)  $\varepsilon = \frac{1}{2}(\nabla \mathbf{u} + \nabla \mathbf{u}^t) = \frac{1}{2}(\mathbf{F} + \mathbf{F}^t) - \mathbf{I}$ 

• Written out:

$$\varepsilon = \frac{1}{2} \begin{pmatrix} 2\partial_x u & \partial_y u + \partial_x v & \partial_z u + \partial_x w \\ \partial_x v + \partial_y u & 2\partial_y v & \partial_z v + \partial_y w \\ \partial_x w + \partial_z u & \partial_y w + \partial_z v & 2\partial_z w \end{pmatrix}$$

Notation
$$\mathbf{u}(\mathbf{x}) = \begin{pmatrix} u(x, y, z) \\ v(x, y, z) \\ w(x, y, z) \end{pmatrix}$$

#### **3D Linear Strain**

• Linear Cauchy strain

$$\mathcal{E} = \frac{1}{2} \begin{pmatrix} 2\partial_x u & \partial_y u + \partial_x v & \partial_z u + \partial_x w \\ \partial_x v + \partial_y u & 2\partial_y v & \partial_z v + \partial_y w \\ \partial_x w + \partial_z u & \partial_y w + \partial_z v & 2\partial_z w \end{pmatrix} =: \begin{pmatrix} \varepsilon_x & \gamma_{xy} & \gamma_{xz} \\ \gamma_{xy} & \varepsilon_y & \gamma_{yz} \\ \gamma_{xz} & \gamma_{yz} & \varepsilon_z \end{pmatrix}$$

 $\mathcal{E}_i$ : normal strains  $\mathcal{Y}_i$ : shear strains

• Geometric interpretation



### Cauchy vs. Green strain

Nonlinear Green strain is rotation-invariant

• Apply incremental rotation **R** to given deformation **F** to obtain  $\mathbf{F}' = \mathbf{RF}$ 

• Then 
$$\mathbf{E}' = \frac{1}{2} (\mathbf{F}'^T \mathbf{F}' - \mathbf{I}) = \mathbf{E}$$

◆ Linear Cauchy strain
 is not rotation-invariant
 ε' = <sup>1</sup>/<sub>2</sub> (**F**' + **F**'<sup>t</sup>) ≠ ε
 → artifacts for larger rotations



#### Mass Spring vs Continuum **Mechanics**

Mass spring systems:

- 1. Measure of Deformation
- **2.** Material Model k

$$\left(\frac{l}{l_0} - 1\right)$$

- 3. Deformation Energy  $W = \frac{1}{2}ke^2$ 4. Internal Forces  $f_{int} = -\frac{\partial W}{\partial x}$

Continuum Mechanics:

**1.** Measure of Deformation: Green or Cauchy strain

2. Material Model

# Material Model: linear isotropic material

- Material model links strain to energy (and stress)
- Linear isotropic material (*generalized Hooke's law*)
  - Energy density  $\Psi = \frac{1}{2}\lambda tr(\boldsymbol{\varepsilon})^2 + \mu tr(\boldsymbol{\varepsilon}^2)$
  - Lame parameters  $\lambda$  and  $\mu$  are material constants related to Poisson Ratio and Young's modulus
- Interpretation
  - $-\operatorname{tr}(\boldsymbol{\varepsilon}^2) = \|\boldsymbol{\varepsilon}\|_F^2$  penalizes all strain components equally
  - $\operatorname{tr}(\boldsymbol{\varepsilon})^2$  penalizes dilatations, i.e., volume changes

#### Volumetric Strain (dilatation, hydrostatic strain)

- Consider a cube with side length a
- For a given deformation ε, the volumetric strain is
  - $\Delta V/V0 = (a(1+\varepsilon_{11}) \cdot a(1+\varepsilon_{22}) \cdot a(1+\varepsilon_{33}) a^3)/a^3$  $= (\varepsilon_{11} + \varepsilon_{22} + \varepsilon_{33}) + O(\varepsilon^2) \approx \operatorname{tr}(\varepsilon)$



http://en.wikipedia.org/wiki/Infinitesimal\_strain\_theory

#### Linear isotropic material

Energy density: 
$$\Psi = \frac{1}{2}\lambda tr(\boldsymbol{\varepsilon})^2 + \mu tr(\boldsymbol{\varepsilon}^2)$$

• Problem: Cauchy strain is not invariant under rotations  $\rightarrow$  artifacts for rotations

- Solutions:
  - Corotational elasticity
  - Nonlinear elasticity



# Material Model: non-linear isotropic model

- ♦ Replace Cauchy strain with Green strain → St.
  Venant-Kirchhoff material (StVK)
- Energy density:  $\Psi_{StVK} = \frac{1}{2}\lambda tr(\mathbf{E})^2 + \mu tr(\mathbf{E}^2)$
- Rotation invariant!

#### **Problems with StVK**

#### StVK softens under compression

$$\Psi_{StVK} = \frac{1}{2}\lambda tr(\mathbf{E})^2 + \mu tr(\mathbf{E}^2)$$





#### **Advanced nonlinear materials**

- Green Strain  $\mathbf{E} = \frac{1}{2}(\mathbf{F}^t\mathbf{F} \mathbf{I}) = \frac{1}{2}(\mathbf{C} \mathbf{I})$
- ◆ Split into *deviatoric (i.e. shape changing/distortion)* and *volumetric (dilation, volume changing)* deformations
  Volumetric: J = det(F) Deviatoric: Ĉ = det(F)<sup>-2/3</sup> C

Neo-Hookean material:

$$\Psi_{NH} = \frac{\mu}{2} \operatorname{tr} \left( \widehat{\mathbf{C}} - \mathbf{I} \right) - \mu \ln(J) + \frac{\lambda}{2} \ln(J)^2$$



#### **Different Models**



#### St. Venant-Kirchoff

#### Neo-Hookean

Linear

#### Mass Spring vs Continuum Mechanics

Mass spring systems:

- 1. Measure of Deformation
- **2.** Material Model k
- 3. Deformation Energy 4. Internal Forces  $f_{int} = -\frac{\partial W}{\partial r}$

$$\left(\frac{l}{l_0} - 1\right)$$

 $\langle 1 \rangle$ 

$$W = \frac{1}{2}ke^2$$

Continuum Mechanics:

 Measure of Deformation: Green or Cauchy strain
 Material Model: linear, StVK, Neo-Hookean, etc
 From Energy Density to Deformation Energy: Finite Element Discretization

#### **Finite Element Discretization**

Divide domain into discrete elements, e.g.,

tetrahedra



- Explicitly store displacement values at nodes  $(\mathbf{x}_i)$ .
- Displacement field everywhere else obtained through interpolation:  $\mathbf{x}(X) = \sum N_i(X)\mathbf{x}_i$
- Deformation Gradient:  $\mathbf{F} = \frac{\partial \mathbf{x}(X)}{\partial X} = \sum_{i} \mathbf{x}_{i} \left(\frac{\partial N_{i}}{\partial X}\right)^{t}$

#### **Basis Functions**

- Basis functions (aka shape functions)  $N_i(X_j): \mathbb{R}^3 \to \mathbb{R}$
- Satisfy delta-property:  $N_i(X_j) = \delta_{ij}$
- Simplest choice: linear basis functions

$$N_i(\bar{x}, \bar{y}, \bar{z}) = a_i \bar{x} + b_i \bar{y} + c_i \bar{z} + d_i$$

• Compute  $N_i$  (and  $\frac{\partial N_i}{\partial X}$ ) through

$$egin{pmatrix} x_1 & y_1 & z_1 & 1 \ x_2 & y_2 & z_2 & 1 \ x_3 & y_3 & z_3 & 1 \ x_4 & y_4 & z_4 & 1 \ \end{pmatrix} egin{pmatrix} a_i \ b_i \ c_i \ d_i \ \end{pmatrix} \ = \ egin{pmatrix} \delta_{1i} \ \delta_{2i} \ \delta_{3i} \ \delta_{4i} \ \end{pmatrix}$$



### **Constant Strain Elements** (Linear Basis Functions)

- Displacement field is continuous in space
- Deformation Gradient, strain, stress are not
  - Constant Strain per element
- Deformation Gradient can be computed as
  - $\mathbf{F} = eE^{-1}$  where e and E are matrices whose columns are edge vectors in undeformed and deformed configurations

#### Constant Strain Elements: From energy density to deformation energy

• Integrate energy density over the entire element:  $W^e = \int_{\Omega} \Psi(\mathbf{F})$ 

- If basis functions are linear:
  - $\mathbf{F}$  is linear in  $\mathbf{x}_i$
  - **F** is constant throughout element:  $W^{e} = \int_{\Omega} \Psi(\mathbf{F}) = \Psi(\mathbf{F}) \cdot \overline{V}^{e}$
## Mass Spring vs Finite Element Method

Mass spring systems:

- **1.** Measure of Deformation  $\left(\frac{l}{l_0}-1\right)$
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$$W = \frac{1}{2}ke^2$$

- Continuum Mechanics:
  - 1. Measure of Deformation: Green or Cauchy strain 2. Material Model: linear, StVK, Neo-Hookean, etc 3. Deformation Energy: integrate over elements 4. Internal Forces:  $f_{int} = -\frac{\partial W}{\partial r}$

- Discretize into elements (triangles/tetraderons, etc)
- For each element
  - Compute deformation gradient  $\mathbf{F} = eE^{-1}$
  - Use material model to define energy density  $\Psi(\mathbf{F})$
  - Integrate over elements to compute energy: W
  - Compute nodal forces as:  $f_{int} = -\frac{\partial W}{\partial x}$

St. Venant-Kirchhoff material

Neohookean elasticity

$$\mathbf{E} = \frac{1}{2} (\mathbf{F}^{\mathsf{T}} \mathbf{F} - \mathbf{I}) \qquad \qquad \mathbf{I}_1 = \|\mathbf{F}\|_F^2, \quad \mathbf{J} = \det \mathbf{F}$$
$$\Psi = \mu \|\mathbf{E}\|_F + \frac{\lambda}{2} \operatorname{tr}^2(\mathbf{E}) \quad \Psi = \frac{\mu}{2} (\mathbf{I}_1 - 3) - \mu \log(\mathbf{J}) + \frac{\lambda}{2} \log^2(\mathbf{J})$$

Area/volume of element  
$$f = -\frac{\partial W}{\partial x} = -V \frac{\partial \Psi}{\partial F} \frac{\partial F}{\partial x}$$

First Piola-Kirchhoff stress tensor P

St. Venant-Kirchhoff material

Neohookean elasticity

$$\mathbf{E} = \frac{1}{2} (\mathbf{F}^{\mathsf{T}} \mathbf{F} - \mathbf{I}) \qquad \mathbf{I}_{1} = \|\mathbf{F}\|_{\mathsf{F}}^{2}, \quad \mathbf{J} = \det \mathbf{F}$$

$$\Psi = \mu \|\mathbf{E}\|_{\mathsf{F}} + \frac{\lambda}{2} \operatorname{tr}^{2}(\mathbf{E}) \qquad \Psi = \frac{\mu}{2} (\mathbf{I}_{1} - 3) - \mu \log(\mathbf{J}) + \frac{\lambda}{2} \log^{2}(\mathbf{J})$$

$$\mathbf{P} = \mathbf{F} [2\mu \mathbf{E} + \lambda \operatorname{tr}(\mathbf{E})\mathbf{I}] \qquad \mathbf{P} = \mu (\mathbf{F} - \mathbf{F}^{-\mathsf{T}}) + \lambda \log(\mathbf{J})\mathbf{F}^{-\mathsf{T}}$$
Area/volume of element
$$f = -\frac{\partial W}{\partial \mathbf{x}} = -V \frac{\partial \Psi}{\partial \mathbf{F}} \frac{\partial \mathbf{F}}{\partial \mathbf{x}}$$

First Piola-Kirchhoff stress tensor **P** 

St. Venant-Kirchhoff material

Neohookean elasticity

$$\begin{split} \mathbf{E} &= \frac{1}{2} (\mathbf{F}^{\mathsf{T}} \mathbf{F} - \mathbf{I}) & I_1 = \|\mathbf{F}\|_{\mathsf{F}}^2, \ J = \det \mathbf{F} \\ \Psi &= \mu \|\mathbf{E}\|_{\mathsf{F}} + \frac{\lambda}{2} \mathrm{tr}^2(\mathbf{E}) & \Psi = \frac{\mu}{2} (I_1 - 3) - \mu \log(J) + \frac{\lambda}{2} \log^2(J) \\ \mathbf{P} &= \mathbf{F} \left[ 2\mu \mathbf{E} + \lambda \mathrm{tr}(\mathbf{E}) \mathbf{I} \right] & \mathbf{P} = \mu (\mathbf{F} - \mathbf{F}^{-\mathsf{T}}) + \lambda \log(J) \mathbf{F}^{-\mathsf{T}} \end{split}$$

For a tetrahedron, this works out to:

$$[f_1 \ f_2 \ f_3] = -VPE^{-T}; f_4 = -f_1 - f_2 - f_3$$

Additional reading: http://www.femdefo.org/

#### **Material Parameters**

St. Venant-Kirchhoff material

Neohookean elasticity

$$\begin{split} \mathbf{E} &= \frac{1}{2} (\mathbf{F}^{\mathsf{T}} \mathbf{F} - \mathbf{I}) & \mathbf{I}_1 = \|\mathbf{F}\|_F^2, \ \mathbf{J} = \det \mathbf{F} \\ \Psi &= \mu \|\mathbf{E}\|_F + \frac{\lambda}{2} \mathrm{tr}^2 (\mathbf{E}) & \Psi = \frac{\mu}{2} (\mathbf{I}_1 - 3) - \mu \log(\mathbf{J}) + \frac{\lambda}{2} \log^2(\mathbf{J}) \end{split}$$

Lame parameters  $\lambda$  and  $\mu$  are material constants related to the fundamental physical parameters: Poisson's Ratio and Young's modulus (http://en.wikipedia.org/wiki/Lamé\_parameters)

## Young's Modulus and Poisson Ratio

Lame parameters  $\lambda$  and  $\mu$  are material constants related to the fundamental physical parameters: Poisson's Ratio and Young's modulus (http://en.wikipedia.org/wiki/Lamé\_parameters)



Young's modulus (E), measure of stiffness Poisson's ratio (v), relative transverse to axial deformation

Stiffness is pretty intuitive







Poisson's Ratio controls volume preservation



Poisson's Ratio controls volume preservation



#### Poisson's Ratio controls volume preservation



#### Poisson's Ratio controls volume preservation





OBJECT

PR = 0.5

#### Poisson's ratio is between -1 and 0.5



#### **Negative Poisson's Ratio**



#### Measurement

## Where do material parameters come from?

#### Simple Measurement: Stiffness



What's the Force (Stress)? What's the Deformation (Strain)?











How do we get the stiffness ?



How do we get the stiffness ?



How do we get the stiffness ?







Compute changes in width and height





#### **Measurement Devices**





## Simulating Elastic Materials with CM+FEM

 You now have all the mathematical tools you need

# Suppose you want to simulate this...



## **Plastic and Elastic Materials**

#### Elastic Materials

- Objects return to their original shape in the absence of other forces
- Plastic Deformations:
  - Object does not always return to its original shape

#### **Example: Crushing a Coke Can**





#### **Old Reference State**

**New Reference State** 

#### **Example: Crushing a van**



## **A Simple Model For Plasticity**

- Recall our model for strain:  $\frac{1}{2} (\mathbf{F}^T \mathbf{F} \mathbf{I})$
- Let's consider how to encode a change of reference shape into this metric
  - Changing undeformed mesh is not easy!
- $\blacklozenge$  We want to exchange F with  $_p^w {\bf F}$ , a deformation gradient that takes into account the new shape of our object



**New Reference State** 

## **Continuum Mechanics: Deformation**

 deformation gradient maps undeformed vectors (local) to deformed (world) vectors


deformation gradient maps undeformed vectors (reference) to deformed (world) vectors



 F transforms a vector from Reference space to World Space



#### Introduce a new space







 $\begin{array}{ccc} \bullet & \text{Our goal is to use} & {}^w_p \mathbf{F} & \text{but we only have} \\ \text{access to} & {}^w_r \mathbf{F} & \\ & {}^w_p \mathbf{F} = {}^w_r \mathbf{F}{}^p_r \mathbf{F}{}^{-1} \\ & {}^w_p \mathbf{F} = {}^w_r \mathbf{F}{}^p_r \mathbf{F}{}^{-1} \end{array}$ 

• Keep an estimate of  ${}^{p}_{r}\mathbf{F}^{-1}$  per element, built incrementally

## How to Compute the Plastic Deformation Gradient

- We compute the strain/stress for each element during simulation
- When it gets above a certain threshold store F as  ${}^p_{rr}\mathbf{F}$



## How to Compute the Plastic Deformation Gradient

- We compute the strain/stress for each element during simulation
- When it gets above a certain threshold store F as  ${}^p_{rr}\mathbf{F}$



#### How to Compute the Plastic Deformation Gradient

- Each subsequent simulation step uses
  - $_{p}^{w}\mathbf{F} =_{r}^{w}\mathbf{F}_{r}^{p}\mathbf{F}^{-1}$



# So now you too can simulate this...



#### **Questions?**