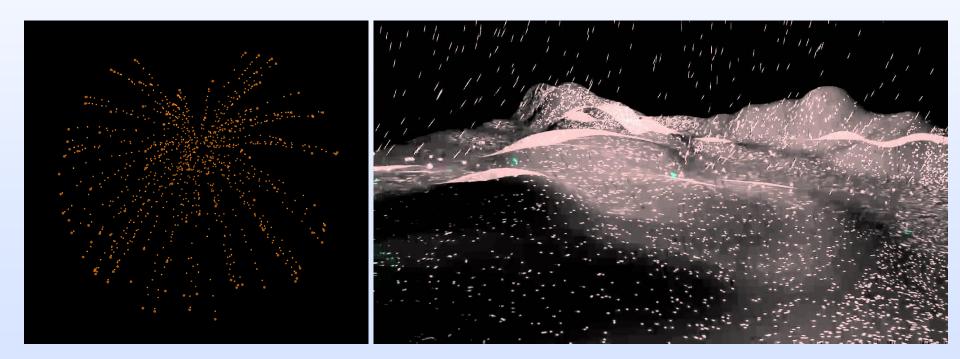
## Mass-Spring Systems



## What we've seen so far

## Individual particles moving under the influence of various types of forces

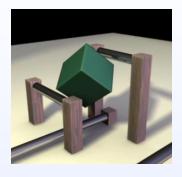


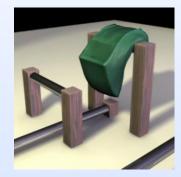
### Material diversity

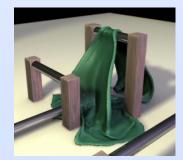
- Deformable objects
  - deform under applied forces
  - resist deformation

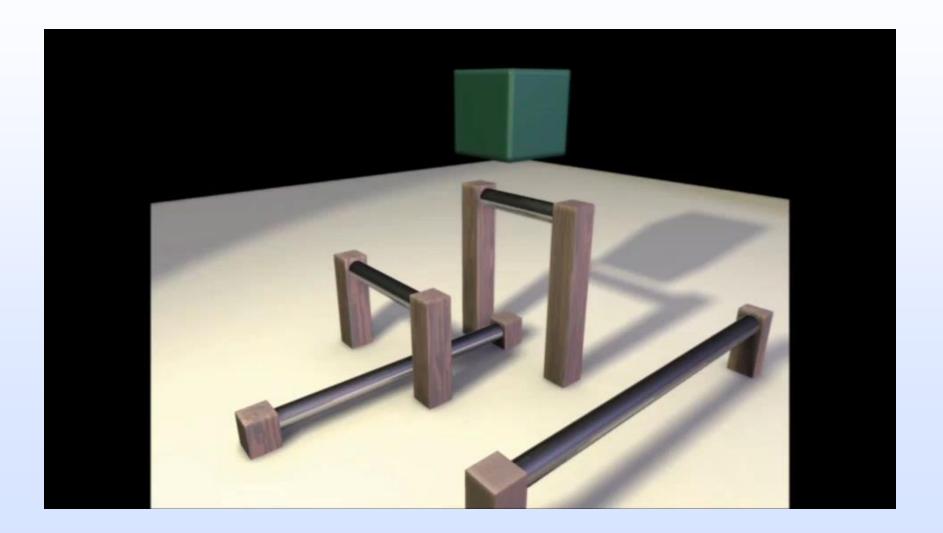
#### Common material properties

- Elastic: deformations are reversible
- Viscous: amplitude of oscillations is reduced
- **Plastic**: irreversible deformations
- Different combinations, different extents



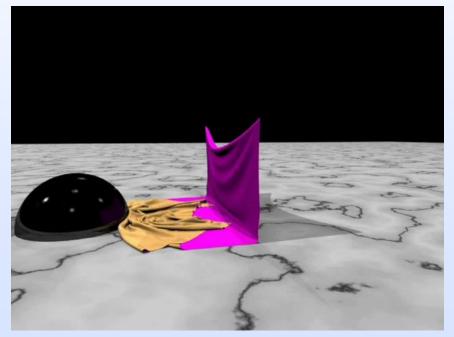






### Model diversity

#### **Cloth simulation**



Bridson et al., 2002

#### Hair simulation



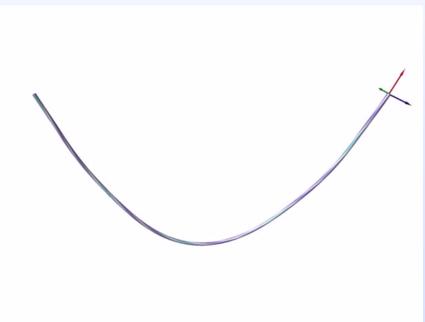
Selle et al., 2008

### Model diversity

#### **Thin Shells**



Rods



[Bridson et al. '03]

#### [Bergou et al. '08]

#### **Diversity of phenomena**

#### Adaptive Tearing and Cracking of Thin Sheets

Tobias Pfaff Rahul Narain Juan Miguel de Joya James F. O'Brien

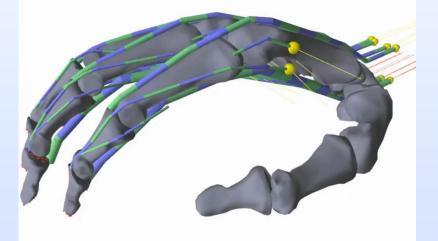
UC Berkeley



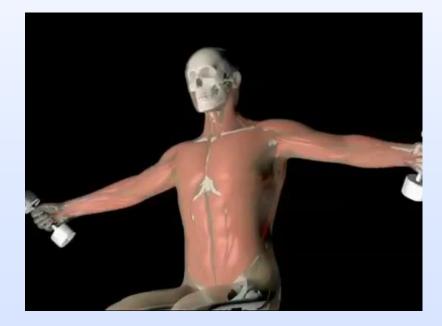
#### Beyond passive objects

#### **Muscles & biological tissues**

Forming a fist

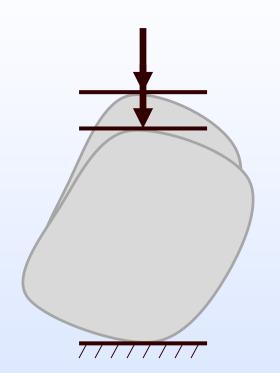


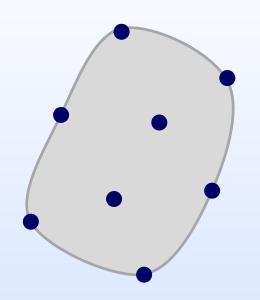
[Sachdeva et al. '15]



[Sifakis et al. '15]

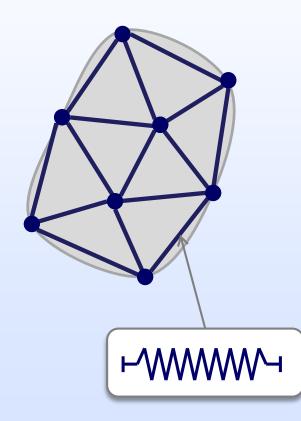
#### Modeling complex phenomena

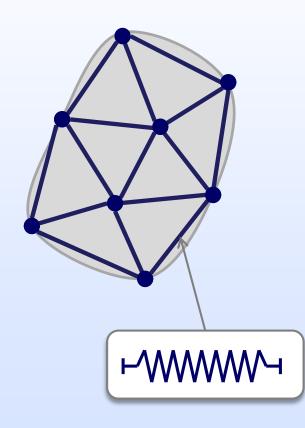




Spatial discretization: sample object with mass points

- ♦ Total mass of object: M
- Number of mass points: p
- Mass of each point: m=M/p (uniform distribution)
- Each point is a particle, just like before. It has
- Mass  $m_i$
- Position  $\mathbf{x}_i(t)$
- Velocity  $\mathbf{v}_i(t)$





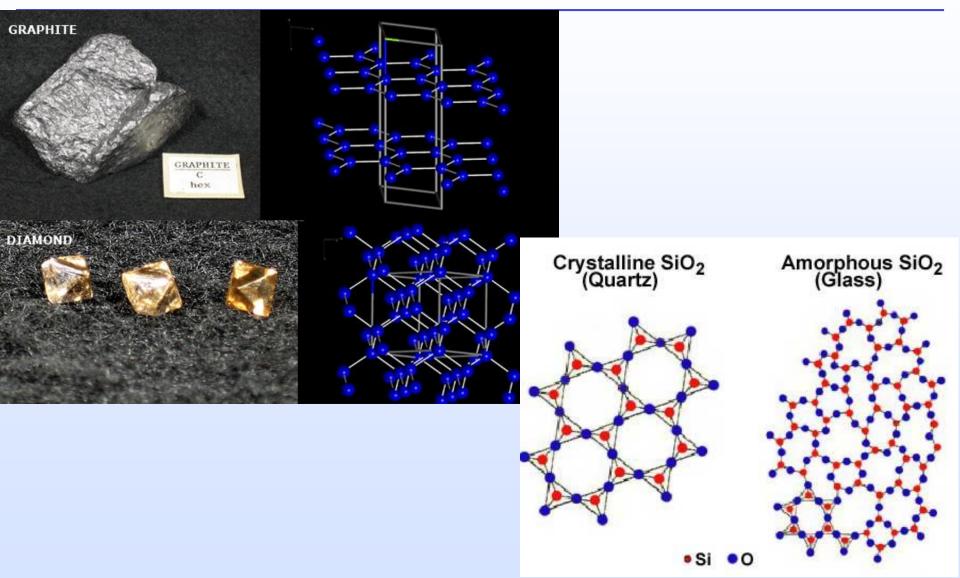
#### **External forces**

- Gravity
- Contact & Collision Forces

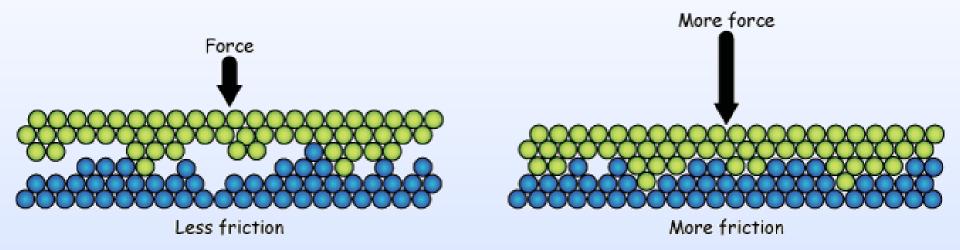
#### **Internal forces**

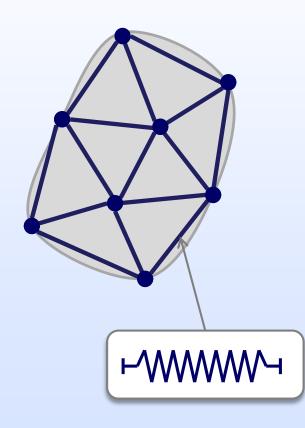
- Elastic spring forces
- Viscous damping forces
- Should always sum up to zero!

# "Mass-spring" systems in the wild



# "Mass-spring" systems in the wild



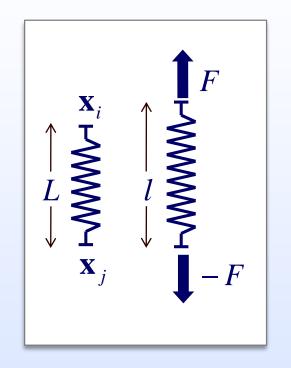


#### **External forces**

- Gravity
- Contact & Collision Forces

#### **Internal forces**

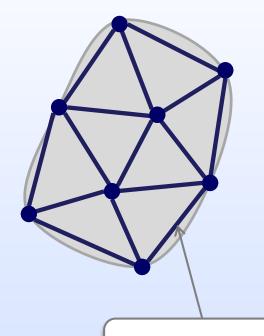
- Elastic spring forces
- Viscous damping forces
- Should always sum up to zero!



Zero length spring  $F_{spring} = -k(x_i - x_j)$ Non-zero length spring  $F_{spring} = -k\left(\frac{|x_i - x_j|}{L} - 1\right)\frac{x_i - x_j}{|x_i - x_j|}$ 

Initial spring lengthLSpring stiffnessk

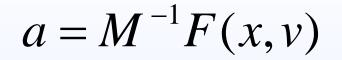
16

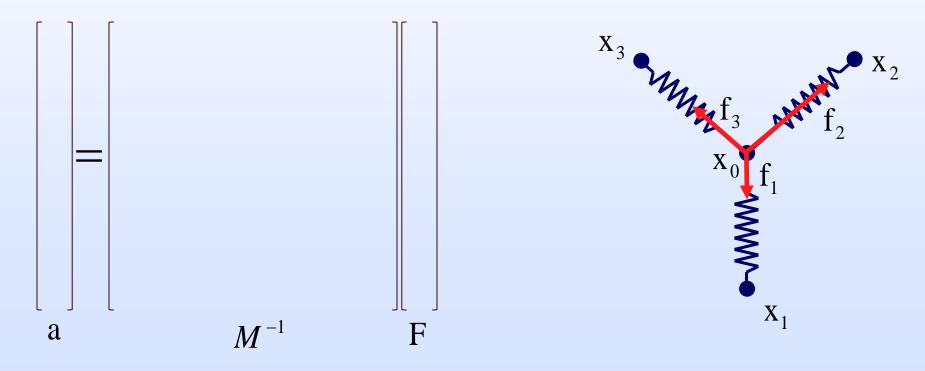


н-////////--

Modularity is key here. If you know how to model one spring, you know how to model complex objects!

## **Back to Equations of Motion**





## **Back to Equations of Motion**

 $a = M^{-1}F(x, v)$ 

 $a^{i} = \frac{1}{m^{i}} \sum f^{i}(x^{i}, v^{i}) = \frac{1}{m^{i}} F^{i}(x^{i}, v^{i})$ 

## **Symplectic Euler**

1) Compute & sum up forces:  $F(x_n, v_n)$ 2) Update velocities:  $v_{n+1} = v_n + hM^{-1}F(x_n, v_n)$ 3) Update positions:  $x_{n+1} = x_n + hv_{n+1}$ 

## **Backward Euler**

#### Recall need to solve:

$$x_{n+1} = x_n + hv_{n+1}$$
$$v_{n+1} = v_n + hM^{-1}F(x_{n+1}, v_{n+1})$$

• which boils down to solving systems of linear equations:

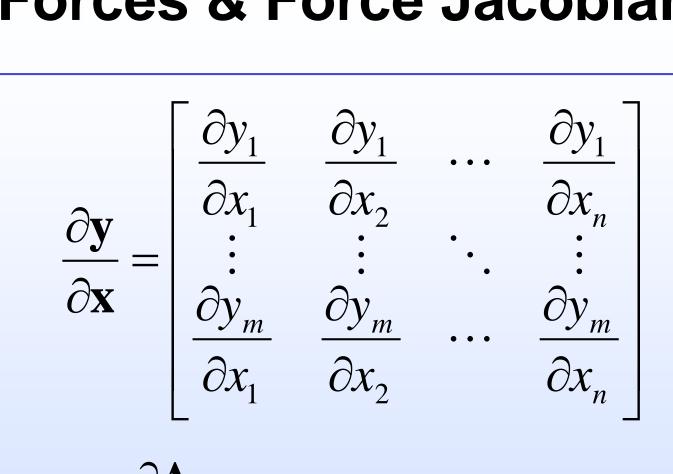
$$\underbrace{\begin{pmatrix} M - h\frac{\partial F}{\partial v} - h^2\frac{\partial F}{\partial x} \\ A & X = \end{pmatrix}}_{A} \Delta v = M\left(v_n - v^k\right) + hF$$

• We've seen what *F* is and how it is computed. What are  $\frac{\partial F}{\partial x}$  and  $\frac{\partial F}{\partial y}$ ?

A bit of tensor calculus overview:

 $\frac{\partial \alpha}{\partial \beta}$ 

derivative of every component of  $\alpha$  with respect to every component of  $\beta$ , where  $\alpha$ and  $\beta$  can be scalars, vectors or matrices



What is  $\frac{\partial A}{\partial x}$  ? A vector of matrices... we will see it soon enough...

- So, what does  $\frac{\partial F}{\partial x}$  represent?
- Much more convenient to think in terms of blocks of the Jacobian:

$$\frac{\partial F_i}{\partial x_j}$$

- - Need to look at derivatives of individual spring forces

#### But first, some useful derivatives

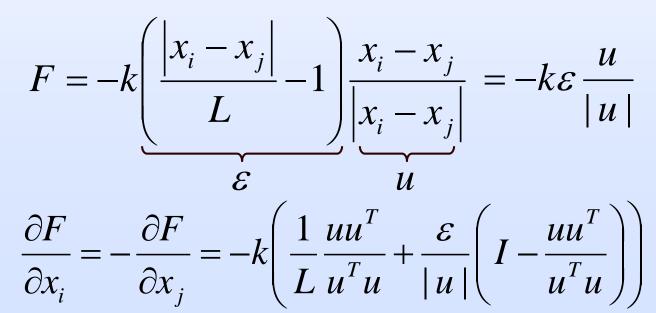
 $\frac{dx^T A x}{2} = x^T (A + A^T)$  $\frac{dA}{\partial x} = 0$  $u = x_i - x_i$  $\frac{du}{\partial x_i} = -\frac{du}{\partial x_i} = I$  $\frac{dx}{\partial x} = I$  $\frac{dx^T x}{\partial x} = 2x^T$  $\frac{d|u|}{\partial x_i} = \frac{u}{|u|}$  $\frac{dAx}{\partial x} = A$  $\frac{da^T x}{\partial x} = \frac{dx^T a}{\partial x} = a^T$  $\frac{du \times v}{\partial x_i} = -v_{\times}$  $\frac{dx^T A}{dx^T A} = A^T$  $v_{\star}$  is a skew-symmetric cross product matrix

Zero length spring

$$F = -k(x_i - x_j)$$

$$\frac{\partial F}{\partial x_i} = -\frac{\partial F}{\partial x_j} = -kI$$

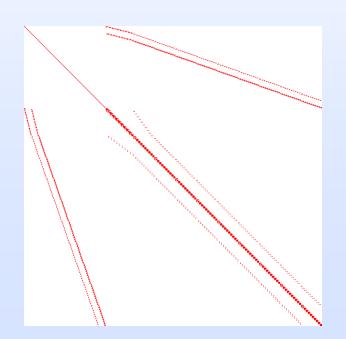
Non-zero Length Spring



- Analytic formulas
- Numerical Approach
  - Finite Differences, very useful for prototyping/debugging
- Automatic & Symbolic differentiation
  - e.g. Maple

We now know how to compute individual entries of the Jacobian, but what does it look like globally?

Block *i*,*j* is non-zero only if there is a spring between particles *i* and *j*. In general, connectivity structure is very sparse - most entries are therefore zero!



## **Backward Euler**

 Boils down to solving systems of linear equations:

$$\underbrace{\begin{pmatrix} M - h\frac{\partial F}{\partial v} - h^2\frac{\partial F}{\partial x} \\ A & X = \end{pmatrix}}_{A} \Delta v = M\left(v_n - v^k\right) + hF$$

- Matrix A is large, sparse, symmetric, (sometimes positive definite)
  - these characteristics will inform the choice of algorithm we can/should use to solve the systems of equations